Grand Unification in the Spectral Pati-Salam Model

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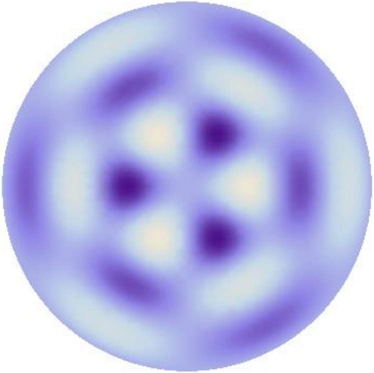
"Can one hear the shape of a drum?" (Kac, 1966)

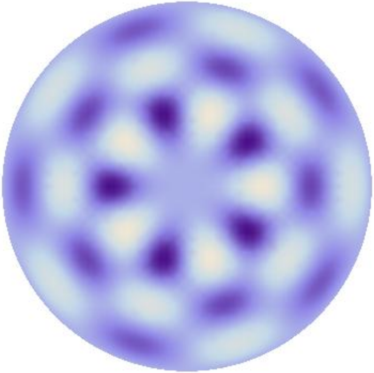
Or, more precisely, given a Riemannian manifold M, does the spectrum of wave numbers k in the Helmholtz equation

$$\boxed{\Delta_M u = k^2 u}$$

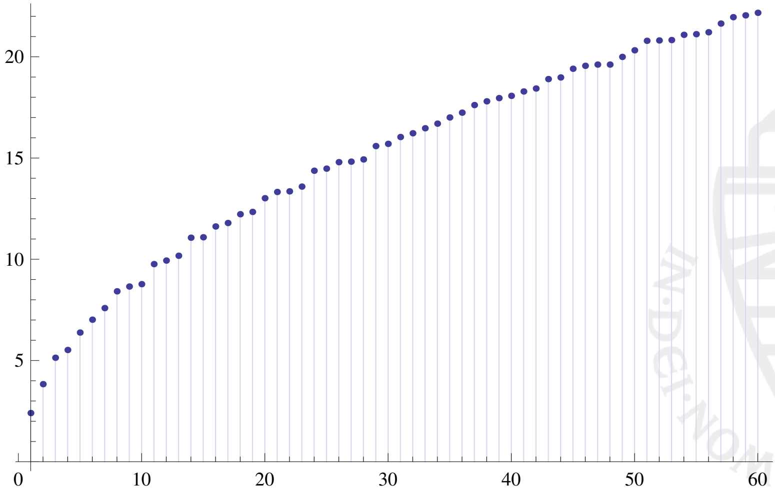
determine the geometry of M?

The disc

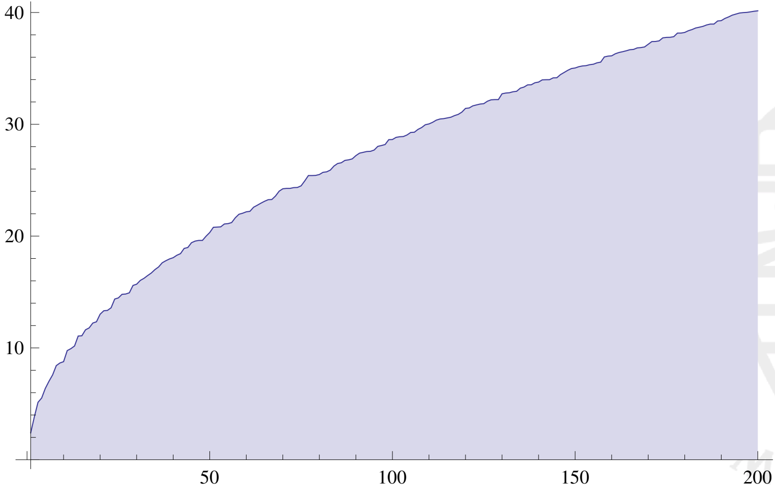




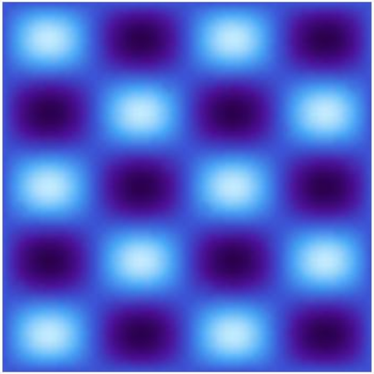
Wave numbers on the disc

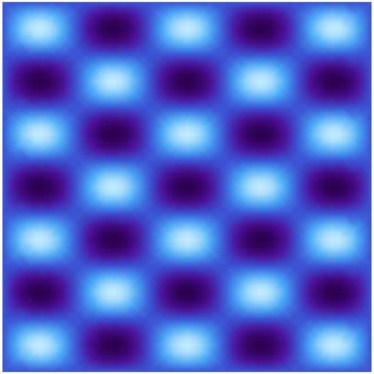


Wave numbers on the disc: high frequencies

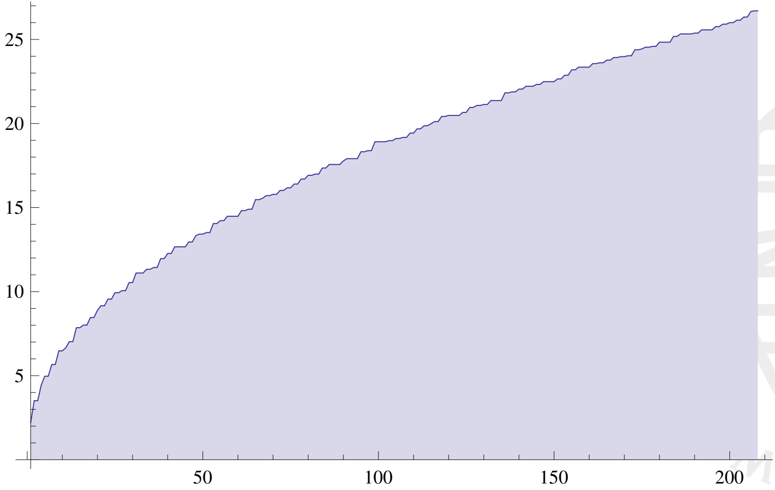


The square



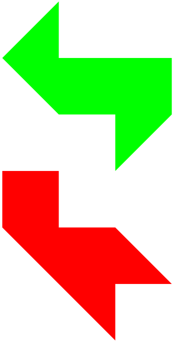


Wave numbers on the square



Isospectral domains

But, there are isospectral domains in $\mathbb{R}^2$:



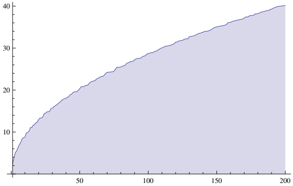
(Gordon, Webb, Wolpert, 1992)

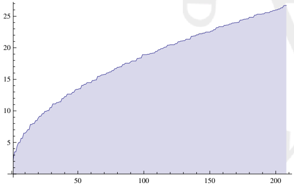
so the answer to Kac's question is no.

Nevertheless, certain information can be extracted from spectrum, such as dimension n of M:

$$N(\Lambda) = \#\text{wave numbers } \leq \Lambda$$
$$\sim \frac{\Omega_n \mathrm{Vol}(M)}{n(2\pi)^n} \Lambda^n$$

For the disc and square this is confirmed by the parabolic shapes ($\sqrt{\Lambda}$):





Recall that k^2 is an eigenvalue of the Laplacian in the Helmholtz equation.

- The Dirac operator is a 'square-root' of the Laplacian, so that its spectrum give the wave numbers k.
- First found by Paul Dirac in flat space, but exists on any Riemannian spin manifold M.
- Let us give some examples.

The circle

- The Laplacian on the circle $\mathbb{S}^1$ is given by

$$\Delta_{\mathbb{S}^1} = -\frac{d^2}{dt^2}; \qquad (t \in [0, 2\pi))$$

- The Dirac operator on the circle is

$$D_{\mathbb{S}^1} = -i\frac{d}{dt}$$

with square $\Delta_{\mathbb{S}^1}$.

- The eigenfunctions of $D_{\mathbb{S}^1}$ are the complex exponential functions

$$e^{int} = \cos nt + i \sin nt$$

with eigenvalue $n \in \mathbb{Z}$.

The 2-dimensional torus

- Consider the two-dimensional torus $\mathbb{T}^2$ parametrized by two angles $t_1, t_2 \in [0, 2\pi)$.
- The Laplacian reads

$$\Delta_{\mathbb{T}^2} = -\frac{\partial^2}{\partial t_1^2} - \frac{\partial^2}{\partial t_2^2}.$$

- At first sight it seems difficult to construct a differential operator that squares to $\Delta_{\mathbb{T}^2}$:

$$\left(a\frac{\partial}{\partial t_1} + b\frac{\partial}{\partial t_2} \right)^2 = a^2\frac{\partial^2}{\partial t_1^2} + 2ab\frac{\partial^2}{\partial t_1 \partial t_2} + b^2\frac{\partial^2}{\partial t_2^2}$$

- This puzzle was solved by Dirac who considered the possibility that a and b be complex *matrices*:

$$a = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}; \qquad b = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix}$$

then $a^2 = b^2 = -1$ and $ab + ba = 0$

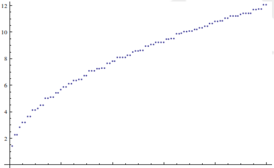
- The Dirac operator on the torus is

$$D_{\mathbb{T}^2} = \begin{pmatrix} 0 & \frac{\partial}{\partial t_1} + i \frac{\partial}{\partial t_2} \\ -\frac{\partial}{\partial t_1} + i \frac{\partial}{\partial t_2} & 0 \end{pmatrix},$$

which satisfies $(D_{\mathbb{T}^2})^2 = -\frac{\partial^2}{\partial t_1^2} - \frac{\partial^2}{\partial t_2^2}$.

- The spectrum of the Dirac operator $D_{\mathbb{T}^2}$ is

$$\left\{ \sqrt{n_1^2 + n_2^2} : n_1, n_2 \in \mathbb{Z} \right\};$$



The 4-dimensional torus

- Consider the 4-torus $\mathbb{T}^4$ parametrized by t_1, t_2, t_3, t_4 and the Laplacian is

$$\Delta_{\mathbb{T}^4} = -\frac{\partial^2}{\partial t_1^2} - \frac{\partial^2}{\partial t_2^2} - \frac{\partial^2}{\partial t_3^2} - \frac{\partial^2}{\partial t_4^2}.$$

- The search for a differential operator that squares to $\Delta_{\mathbb{T}^4}$ again involves matrices, but we also need quaternions:

$$i^2 = j^2 = k^2 = ijk = -1.$$

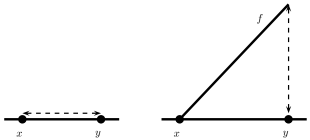
- The Dirac operator on $\mathbb{T}^4$ is

$$D_{\mathbb{T}^4} = \begin{pmatrix} 0 & \frac{\partial}{\partial t_1} + i\frac{\partial}{\partial t_2} + j\frac{\partial}{\partial t_3} + k\frac{\partial}{\partial t_4} \\ -\frac{\partial}{\partial t_1} + i\frac{\partial}{\partial t_2} + j\frac{\partial}{\partial t_3} + k\frac{\partial}{\partial t_4} & 0 \end{pmatrix}$$

- The relations $ij = -ji$, $ik = -ki$, *et cetera* imply that its square coincides with $\Delta_{\mathbb{T}^4}$.

- As said, the geometry of M is not fully determined by spectrum of D_M.
- This can be improved by considering besides D_M also the algebra $C^\infty(M)$ of smooth (coordinate) functions on M
- In fact, the distance function on M is equal to

$$d(x,y) = \sup_{f \in C^\infty(M)} \{|f(x) - f(y)| : \text{ gradient } f \leq 1\}$$



- The gradient of f is given by the commutator $[D_M, f] = D_M f - f D_M$ (e.g. $[D_{\mathbb{S}^1}, f] = -i \dfrac{df}{dt}$)

A fermion in a spacetime background

The combination of coordinates and Dirac operator is of course also central in the description of fermion propagation:

- coordinates on spacetime M:

$$x_\mu \cdot x_\nu(p) = x_\mu(p) x_\nu(p), etc.,$$

  with $\mu, \nu = 1, \ldots, 4$.

- propagation, described by Dirac operator $\not{\partial}_M = i\gamma^\mu \partial_\mu$, acting on wavefunctions ψ:

$$S[\psi] = \int \overline{\psi} \not{\partial}_M \psi \qquad \rightsquigarrow \text{EOM: } \not{\partial}_M \psi = 0.$$

Emerging bosons

Our fermionic starting point induces a bosonic theory:

- "Inner fluctuations" by the coordinates [C 1996]:

$$\partial\!\!\!/_M \rightsquigarrow \partial\!\!\!/_M + \sum_j a_j [\partial\!\!\!/_M, a_j']$$

  for functions a_j, a_j' depending on the coordinates x_μ.

- Then, by the chain rule:

$$\sum_j a_j [\partial\!\!\!/_M, a_j'] = A^\nu \gamma^\mu (\partial_\mu x^\nu) = A^\mu \gamma_\mu$$

  where A^μ is the electromagnetic 4-potential describing the photon.

Moreover, it is possible to derive a bosonic action from the (Euclidean) Dirac operator via the spectral action [CC 1996]:

$$\mathrm{Trace}\, e^{-\not{\partial}_M^2/\Lambda^2} \sim c_4 \Lambda^4 \mathrm{Vol}(M) + c_2 \Lambda^2 \int R\sqrt{g} + c_0 \int (\partial_{[\mu} A_{\nu]})^2 + \cdots$$

for some coefficients $c_4, c_2, \ldots$.

We recognize

- The Einstein-Hilbert action $\int R\sqrt{g}$ for (Euclidean) gravity
- The Lagrangian $\int (\partial_{[\mu} A_{\nu]})^2$ for the electromagnetic field

Replace spacetime by
spacetime $\times$ **noncommutative space**: $M \times F$

- F is considered as finite internal space (Kaluza–Klein like)
- F is described by noncommutative matrices, that play the role of coordinates, just as spacetime is described by $x_\mu(p)$.
- 'Propagation' of particles in F is described by a 'Dirac operator' $\partial\!\!\!/_F$ which is actually simply a hermitian matrix.

Finite **commutative** spaces

- Finite space F

$$F = \quad {}_1\bullet \qquad {}_2\bullet \qquad \cdots\cdots \qquad {}_N\bullet$$

- Coordinate functions on F are given by N-tuples in $\mathbb{C}^N$, and the corresponding algebra $C^\infty(F)$ corresponds to diagonal matrices

$$\begin{pmatrix} f(1) & 0 & \cdots & 0 \\ 0 & f(2) & \cdots & 0 \\ \vdots & & \ddots & \vdots \\ 0 & 0 & \ldots & f(N) \end{pmatrix}$$

- The finite Dirac operator is an arbitrary hermitian matrix D_F, giving rise to a distance function on F as

$$d(p,q) = \sup_{f \in C^\infty(F)} \{ |f(p) - f(q)| : \| [D_F, f] \| \leq 1 \}$$

Example: two-point space

$$F = {}_1\bullet \qquad {}_2\bullet$$

- Then the algebra of smooth functions

$$C^\infty(F) := \left\{ \begin{pmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{pmatrix} \middle| \lambda_1, \lambda_2 \in \mathbb{C} \right\}$$

- A finite Dirac operator is given by

$$D_F = \begin{pmatrix} 0 & \overline{c} \\ c & 0 \end{pmatrix}; \qquad (c \in \mathbb{C})$$

- The distance formula then becomes

$$d(1,2) = \frac{1}{|c|}$$

Finite **noncommutative** spaces

The geometry of F gets much more interesting if we allow for a *noncommutative* structure at each point of F.

- Instead of diagonal matrices, we consider block diagonal matrices

$$A = \begin{pmatrix} a_1 & 0 & \cdots & 0 \\ 0 & a_2 & \cdots & 0 \\ \vdots & & \ddots & \vdots \\ 0 & 0 & \dots & a_N \end{pmatrix},$$

  where the $a_1, a_2, \dots, a_N$ are square matrices of size $n_1, n_2, \dots, n_N$.

- Hence we will consider the matrix algebra

$$A_F := M_{n_1}(\mathbb{C}) \oplus M_{n_2}(\mathbb{C}) \oplus \cdots \oplus M_{n_N}(\mathbb{C})$$

  where $\mathbb{C}$ can be replaced by $\mathbb{R}$ or $\mathbb{H}$.

- A finite Dirac operator is still given by a hermitian matrix.

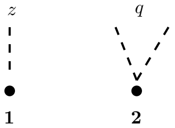
Example: noncommutative two-point space

Coordinates on F are elements in $\mathbb{C} \oplus \mathbb{H}$

- A complex number z
- A quaternion $q = q_0 + iq_k\sigma^k$; in terms of Pauli matrices:

$$\sigma^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma^2 = \begin{pmatrix} 0 & i \\ -i & 0 \end{pmatrix}, \quad \sigma^3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

It describes a two-point space, with internal structure:



Gauge group is given by unitaries: $U(1) \times SU(2)$.

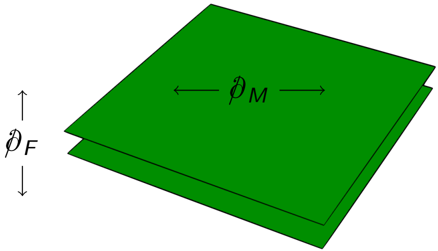
'Dirac operator'

$$\not{\partial}_F = \begin{pmatrix} 0 & \overline{c} & 0 \\ c & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

- "Inner fluctuations" can be defined as before but now yield:

$$\sum_j \begin{pmatrix} z_j & 0 \\ 0 & q_j \end{pmatrix} \left[\not{\partial}_F, \begin{pmatrix} z_j' & 0 \\ 0 & q_j' \end{pmatrix} \right] =: \begin{pmatrix} 0 & \overline{c}\overline{\phi_1} & \overline{c}\overline{\phi_2} \\ c\phi_1 & 0 & 0 \\ c\phi_2 & 0 & 0 \end{pmatrix}$$

Almost-commutative spacetimes



We combine this mild (matrix) noncommutativity with spacetime:

- coordinates of the almost-commutative spacetime $M \times F$:

$$\hat{x}^\mu(p) = (z^\mu(p), q^\mu(p))$$

  as elements in $\mathbb{C} \oplus \mathbb{H}$ (for each μ and each point p of M)
- The combined Dirac operator becomes

$$\not{D}_{M\times F} = \not{D}_M + \gamma_5 \not{D}_F$$

  Note that $\not{D}_{M\times F}^2 = \not{D}_M^2 + \not{D}_F^2$, which will be useful later on.

Inner fluctuations on $M \times F$

So, we describe $M \times F$ by:

$$\hat{x}^\mu = (z^\mu, q^\mu)\,; \qquad \partial\!\!\!/_{M\times F} = \partial\!\!\!/_M + \gamma_5 \partial\!\!\!/_F$$

As before, we consider inner fluctuations of $\partial\!\!\!/_{M\times F}$ by $\hat{x}^\mu(p)$:

- The inner fluctuations of $\partial\!\!\!/_F$ become scalar fields ϕ_1, ϕ_2.
- The inner fluctuations of $\partial\!\!\!/_M$ become matrix-valued:

$$\sum_j a_j [\partial\!\!\!/_M, a_j'] = a_\nu \gamma^\mu (\partial_\mu \hat{x}^\nu) =: \partial\!\!\!/_M + A_\mu \gamma^\mu$$

with A_μ taking values in $\mathbb{C} \oplus \mathbb{H}$:

$$A_\mu = \begin{pmatrix} B_\mu & 0 & 0 \\ 0 & W_\mu^3 & W_\mu^+ \\ 0 & W_\mu^- & -W_\mu^3 \end{pmatrix}$$

corresponding to hypercharge and the W-bosons.

Action functional: electroweak theory

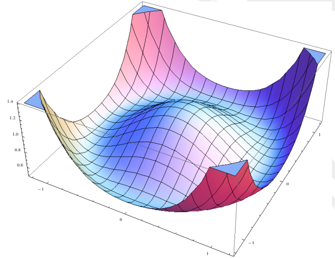
Use $\not{\partial}^2_{M\times F} = \not{\partial}^2_M + \not{\partial}^2_F$ to compute the spectral action

$$\mathrm{Trace}\, e^{-\not{\partial}^2_{M\times F}/\Lambda^2} = \mathrm{Trace}\, e^{-\not{\partial}^2_M/\Lambda^2}\left(1 - \frac{\not{\partial}^2_F}{\Lambda^2} + \frac{1}{2}\frac{\not{\partial}^4_F}{\Lambda^4} - \cdots\right)$$

$$\sim \left(c_4\Lambda^4 \mathrm{Vol}(M) + c_2\Lambda^2 \int R\sqrt{g} + c_0 \int F_{\mu\nu}F^{\mu\nu}\right)\left(1 - \frac{|\phi|^2}{\Lambda^2} + \frac{|\phi|^4}{2\Lambda^4}\right) + \cdots$$

We now recognize in terms of the field-strength $F_{\mu\nu}$ for A_μ:

- The Yang–Mills term $F_{\mu\nu}F^{\mu\nu}$ for hypercharge and W-boson
- The Higgs potential $-c_4\Lambda^2|\phi|^2 + \frac{1}{2}c_4|\phi|^4$



Standard Model as an almost-commutative spacetime

Describe $M \times F_{SM}$ by [CCM 2007]

- Coordinates: $\hat{x}^\mu(p) \in \mathbb{C} \oplus \mathbb{H} \oplus M_3(\mathbb{C})$ (with unimodular unitaries $U(1)_Y \times SU(2)_L \times SU(3)$).
- Dirac operator $\not{\partial}_{M\times F} = \not{\partial}_M + \gamma_5 \not{\partial}_F$ where

$$\not{\partial}_F = \begin{pmatrix} S & T^* \\ T & \overline{S} \end{pmatrix}$$

is a 96×96-dimensional hermitian matrix where 96 is:

$$3 \times 2 \times \left(\underline{2} \otimes \underline{1} + \underline{1} \otimes \underline{1} + \underline{1} \otimes \underline{1} + \underline{2} \otimes \underline{3} + \underline{1} \otimes \underline{3} + \underline{1} \otimes \underline{3} \right)$$

3: families; 2: anti-particles; $\underline{2} \otimes \underline{1}$: (ν_L, e_L); $\underline{1} \otimes \underline{1}$: ν_R; $\underline{1} \otimes \underline{1}$: e_R; $\underline{2} \otimes \underline{3}$: (u_L, d_L); $\underline{1} \otimes \underline{3}$: u_R; $\underline{1} \otimes \underline{3}$: d_R

The Dirac operator on F_{SM}

$$D_F = \begin{pmatrix} S & T^* \\ T & \overline{S} \end{pmatrix}$$

- The operator S is given by

$$S_l := \begin{pmatrix} 0 & 0 & Y_\nu & 0 \\ 0 & 0 & 0 & Y_e \\ Y_\nu^* & 0 & 0 & 0 \\ 0 & Y_e^* & 0 & 0 \end{pmatrix}, \quad S_q \otimes 1_3 = \begin{pmatrix} 0 & 0 & Y_u & 0 \\ 0 & 0 & 0 & Y_d \\ Y_u^* & 0 & 0 & 0 \\ 0 & Y_d^* & 0 & 0 \end{pmatrix} \otimes 1_3,$$

  where Y_ν, Y_e, Y_u and Y_d are 3×3 mass matrices acting on the three generations.

- The symmetric operator T only acts on the right-handed (anti)neutrinos, $T\nu_R = Y_R \overline{\nu_R}$ for a 3×3 symmetric Majorana mass matrix Y_R, and $Tf = 0$ for all other fermions $f \neq \nu_R$.

Inner fluctuations

Just as before, we find

- Inner fluctuations of $\partial\!\!\!/_M$ give a matrix

$$A_\mu = \begin{pmatrix} B_\mu & 0 & 0 & 0 \\ 0 & W_\mu^3 & W_\mu^+ & 0 \\ 0 & W_\mu^- & -W_\mu^3 & 0 \\ 0 & 0 & 0 & (G_\mu^a) \end{pmatrix}$$

  corresponding to hypercharge, weak and strong interaction.

- Inner fluctuations of $\partial\!\!\!/_F$ give

$$\begin{pmatrix} Y_\nu & 0 \\ 0 & Y_e \end{pmatrix} \rightsquigarrow \begin{pmatrix} Y_\nu \phi_1 & -Y_e \overline{\phi}_2 \\ Y_\nu \phi_2 & Y_e \overline{\phi}_1 \end{pmatrix}$$

  corresponding to SM-Higgs field. Similarly for Y_u, Y_d.

Dynamics and interactions

If we reconsider the spectral action:

$$\mathrm{Trace}\, e^{-\not{\partial}^2_{M\times F}/\Lambda^2} \sim \left(c_4 \Lambda^4 \mathrm{Vol}(M) + c_0 \int F_{\mu\nu} F^{\mu\nu} \right) \left(1 - \frac{|\phi|^2}{\Lambda^2} + \frac{|\phi|^4}{2\Lambda^4} \right) + \cdot$$

we observe [CCM 2007]:

- The coupling constants of hypercharge, weak and strong interaction are expressed in terms of the single constant c_0 which implies

$$\boxed{g_3^2 = g_2^2 = \frac{5}{3} g_1^2}$$

  In other words, there should be grand unification.

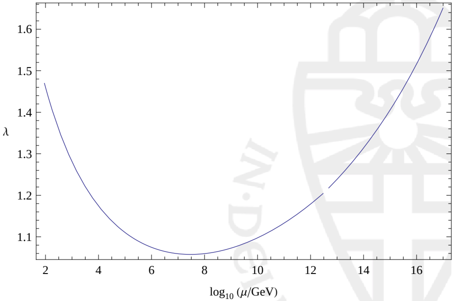
- Moreover, the quartic Higgs coupling λ is related via

$$\lambda \approx 24 \frac{3 + \rho^4}{(3 + \rho^2)^2} g_2^2; \qquad \rho = \frac{m_\nu}{m_{\mathrm{top}}}$$

This can be used to derive predictions as follows:

- Interpret the spectral action as an effective field theory at $\Lambda_{\mathrm{GUT}} \approx 10^{13} - 10^{16}$ GeV.
- Run the quartic coupling constant λ to SM-energies to predict

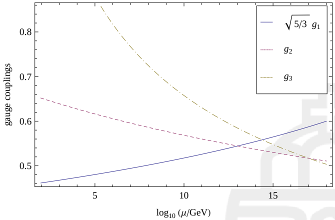
$$m_h^2 = \frac{4\lambda M_W^2}{3g_2^2}$$

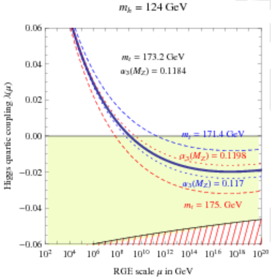


This gives [CCM 2007]

$$\boxed{167 \text{ GeV} \leq m_h \leq 176 \text{ GeV}}$$

1. This prediction is falsified by the now measured value.
2. In the Standard Model there is not the presumed grand unification.
3. There is a problem with the low value of m_h, making the Higgs vacuum un/metastable [Elias-Miro et al. 2011].





- The matrix coordinates of the Standard Model arise naturally as a restriction of the following coordinates

$$\hat{x}^\mu(p) = \left(q_R^\mu(p), q_L^\mu(p), m^\mu(p)\right) \in \mathbb{H}_R \oplus \mathbb{H}_L \oplus M_4(\mathbb{C})$$

corresponding to a Pati–Salam unification:

$$U(1)_Y \times SU(2)_L \times SU(3) \to SU(2)_R \times SU(2)_L \times SU(4)$$

- The 96 fermionic degrees of freedom are structured as

$$\left(\begin{array}{cc|cc} \nu_R & u_{iR} & \nu_L & u_{iL} \\ e_R & d_{iR} & e_L & d_{iL} \end{array} \right) \qquad (i = 1, 2, 3)$$

- Again the finite Dirac operator is a 96×96-dimensional matrix (details in [CCS 2013]).

Inner fluctuations

- Inner fluctuations of $\not{\partial}_M$ now give three gauge bosons:

$$W_R^\mu, \qquad W_L^\mu, \qquad V^\mu$$

  corresponding to $SU(2)_R \times SU(2)_L \times SU(4)$.
- For the inner fluctuations of $\not{\partial}_F$ we distinguish two cases, depending on the initial form of $\not{\partial}_F$:
  - I The Standard Model $\not{\partial}_F = \begin{pmatrix} S & T^* \\ T & \overline{S} \end{pmatrix}$
  - II A more general $\not{\partial}_F$ with zero $\overline{f}_L - f_L$-interactions.

Scalar sector of the spectral Pati–Salam model

Case I For a SM $\not{\partial}_F$, the resulting scalar fields are composite fields, expressed in scalar fields whose representations are:

	$SU(2)_R$	$SU(2)_L$	$SU(4)$
$\phi_{\dot{a}}^{b}$	2	2	1
$\Delta_{\dot{a}I}$	2	1	4
Σ_J^I	1	1	15

Case II For a more general finite Dirac operator, we have fundamental scalar fields:

particle	$SU(2)_R$	$SU(2)_L$	$SU(4)$
$\Sigma_{\dot{a}J}^{bJ}$	2	2	$1+15$
$H_{\dot{a}I\dot{b}J}$	3	1	10
	1	1	6

Action functional

As for the Standard Model, we can compute the spectral action which describes the usual Pati–Salam model with

- unification of the gauge couplings

$$g_R = g_L = g.$$

- A rather involved, fixed scalar potential, still subject to further study

However, independently from the spectral action, we can analyze the running at one loop of the gauge couplings [CCS 2015]:

1. We run the Standard Model gauge couplings up to a presumed PS → SM symmetry breaking scale m_R
2. We take their values as boundary conditions to the Pati–Salam gauge couplings g_R, g_L, g at this scale via

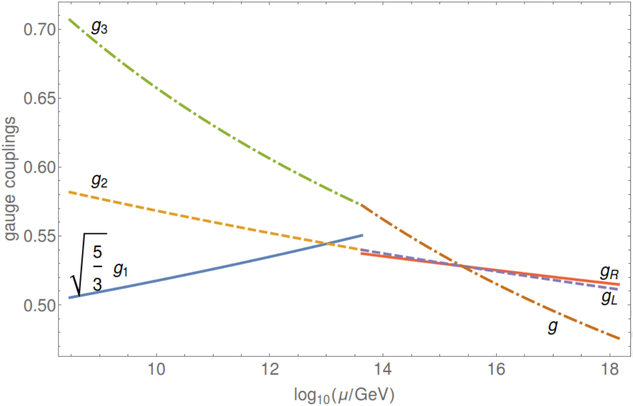
$$\frac{1}{g_1^2} = \frac{2}{3}\frac{1}{g^2} + \frac{1}{g_R^2}, \qquad \frac{1}{g_2^2} = \frac{1}{g_L^2}, \qquad \frac{1}{g_3^2} = \frac{1}{g^2},$$

3. Vary m_R in a search for a unification scale Λ where

$$g_R = g_L = g$$

which is where the spectral action is valid as an effective theory.

For the Standard Model Dirac operator, we have found that with $m_R \approx 4.25 \times 10^{13}$ GeV there is unification at $\Lambda \approx 2.5 \times 10^{15}$ GeV:

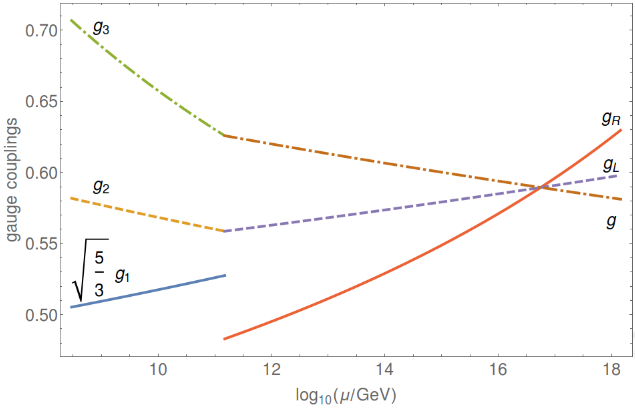


In this case, we can also say something about the scalar particles that remain after SSB:

	$U(1)_Y$	$SU(2)_L$	$SU(3)$
$\begin{pmatrix}\phi_1^0 \\ \phi_1^+\end{pmatrix} = \begin{pmatrix}\phi_{\dot 1}^1 \\ \phi_{\dot 1}^2\end{pmatrix}$	1	2	1
$\begin{pmatrix}\phi_2^- \\ \phi_2^0\end{pmatrix} = \begin{pmatrix}\phi_{\dot 2}^1 \\ \phi_{\dot 2}^2\end{pmatrix}$	-1	2	1
σ	0	1	1
η	$-\frac{2}{3}$	1	3

- It turns out that these scalar fields have a little influence on the running of the SM-gauge couplings (at one loop).
- However, this sector contains the real scalar singlet σ that allowed for a realistic Higgs mass and that stabilizes the Higgs vacuum [CC 2012].

For the more general case, we have found that with $m_R \approx 1.5 \times 10^{11}\ \mathrm{GeV}$ there is unification at $\Lambda \approx 6.3 \times 10^{16}\ \mathrm{GeV}$:



Conclusion

We have arrived at a spectral Pati–Salam model that

- goes beyond the Standard Model
- has a fixed scalar sector once the finite Dirac operator has been fixed (only a few scenarios)
- exhibits grand unification for all of these scenarios (confirmed by [Aydemir–Minic–Sun–Takeuchi 2015])
- the scalar sector has the potential to stabilize the Higgs vacuum and allow for a realistic Higgs mass.

Further reading

A. Chamseddine, A. Connes, WvS.

Beyond the Spectral Standard Model: Emergence of Pati-Salam Unification. *JHEP* 11 (2013) 132. [arXiv:1304.8050]

Grand Unification in the Spectral Pati-Salam Model. *JHEP* 11 (2015) 011. [arXiv:1507.08161]

WvS.

Noncommutative Geometry and Particle Physics. Mathematical Physics Studies, Springer, 2015.

and also: `http://www.noncommutativegeometry.nl`