Factorization of Dirac operators in unbounded KK-theory

Walter van Suijlekom

(joint with Jens Kaad)

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Motivation

Consider the following setup:

• A Riemannian submersion

$$\pi:M\to B$$

of $spin^c$ manifolds M and B

• It is well-known from wrong-way functoriality [CS, 1984] that

$$[M] = \pi! \widehat{\otimes}_{C_0(B)}[B] \tag{*}$$

 Moreover, the Dirac operator on M can be written as a tensor sum:

$$D_{M} = D_{V} \widehat{\otimes} \gamma + 1 \widehat{\otimes}_{\nabla} D_{B} + \widetilde{c}(\Omega)$$

representing the internal KK-product (*) [KvS, 2016].

This talk: can we obtain similar factorization for singular fibrations, that is, if the fiber dimension of $\pi: M \to B$ is allowed to vary?

Key application: toric noncommutative manifolds

- Let \mathbb{T}^n act on \overline{M} by isometries with principal stratum $M \subseteq \overline{M}$.
- Deformation quantization both at the topological [Rie, 1993] and geometric level [CL, 2001]:

$$(C^{\infty}(\overline{M}_{\theta}), L^{2}(\mathscr{E}_{\overline{M}}), D_{\overline{M}})$$

• At the C^* -algebraic level, using a real structure one can identify [vS, 2015] a continuous C^* -bundle $\mathfrak B$ so that

$$C(\overline{M}_{\theta}) = \Gamma(\overline{M}/\mathbb{T}^n, \mathfrak{B})$$

- Gauge theory: Inn(A) acts fiberwise on \mathfrak{B} .
- A tensor sum factorization of the above type means that this bundle picture extends to the geometric level:

$$D_{\overline{M}} = D_V \widehat{\otimes} \gamma + 1 \widehat{\otimes}_{\nabla} D_{M/\mathbb{T}^n} + \widetilde{c}(\Omega)$$

where D_V is a suitably defined family of Dirac operators.

Setup: almost-regular fibrations

We consider singular fibrations in the following sense:

Definition

Let \overline{M} be a closed Riemannian manifold, together with a finite union $P = \bigcup_{j=1}^m P_j$ of compact embedded submanifolds $P_j \subseteq \overline{M}$ each without boundary and of codimension strictly greater than 1. If there exist a Riemannian manifold without boundary B and a proper Riemannian submersion $\pi: M \to B$ with total space $M = \overline{M} \setminus P$, we call the data $(\overline{M}, P, B, \pi)$ an almost-regular fibration.

- We say that $(\overline{M}, P, B, \pi)$ is an almost-regular fibration of spin^c manifolds if both \overline{M} and B are spin^c manifolds.
- Major class of examples given by tori acting on closed
 Riemannian manifolds such that orbit space is connected.

For such almost-regular fibrations $(\overline{M}, P, B, \pi)$ of spin^c manifolds, the factorization can be done [KvS, 2017] by introducing

- a C^* -correspondence X between $C_0(M)$ and $C_0(B)$
- A regular selfadjoint operator

$$D_V : \mathsf{Dom}(D_V) \to X$$

with compact resolvent (family of vertical Dirac operators), defining an unbounded KK-cycle from $C_0(M)$ to $C_0(B)$.

A hermitian connection

$$abla : \mathsf{Dom}(
abla) o X \widehat{\otimes}_{C_0(B)} \Omega^1(B)$$

Main results

Then

- $D_V \widehat{\otimes} \gamma$ and $1 \widehat{\otimes}_{\nabla} D_B$ define symmetric unbounded operators on $X \widehat{\otimes}_{C_0(B)} L^2(\mathscr{E}_B^c)$
- Both $\mathscr{E}_{M}^{c} := \Gamma_{c}^{\infty}(M, \mathcal{S}_{\overline{M}})$ and $\mathscr{E}_{\overline{M}} := \Gamma^{\infty}(\overline{M}, \mathcal{S}_{\overline{M}})$ form a core for the selfadjoint operator $D_{\overline{M}}$ on $\mathsf{Dom}(D_{\overline{M}}) \subset L^{2}(\mathscr{E}_{\overline{M}})$.

Theorem (KvS, 2017)

Up to unitary isomorphism we have the following identity of selfadjoint operators

$$D_V\widehat{\otimes}\gamma + 1\widehat{\otimes}_{\nabla}D_B = D_{\overline{M}} + \widetilde{c}(\Omega).$$

where Ω is the curvature of $\pi: M \to B$.

Unbounded Kasparov modules and fiber bundles

Suppose that

- $\pi: M \to B$ is a proper and surjective submersion
- *M* is equipped with a Riemannian metric in the fiber direction.
- E is a smooth hermitian vector bundle on M and write $\mathscr{E} = \Gamma^{\infty}(M, E)$ and $\mathscr{E}^{c} = \Gamma^{\infty}_{c}(M, E)$.
- ullet g is a first-order vertical differential operator on $\mathscr E$
- \mathscr{D} is vertically elliptic, in the sense that the map $[\mathscr{D}, f](x) : E_x \to E_x$ is invertible whenever $(d_V f)(x)$ is non-trivial.

We define a C^* -correspondence X from $C_0(M)$ to $C_0(B)$ using the hermitian structure on \mathcal{E}^c and integration along the fibers:

$$\langle s,t
angle_X(b):=\int_{F_b}\langle s,t
angle_{\mathscr{E}^c}$$

Moreover, \mathscr{D} gives rise to an unbounded operator $D_0: \mathscr{E}^c \to X$ by setting $D_0(s) = \mathscr{D}(s)$.

Theorem (KvS, 2017)

If the above operator D_0 is symmetric then the closure $D=\overline{D_0}$ is selfadjoint and regular and $m(f)(1+D^2)^{-1/2}$ is a compact operator for all $f\in C_c^\infty(M)$. In other words, (X,D) is an unbounded Kasparov module from $C_0(M)$ to $C_0(B)$.

Riemannian submersions

• Let M and B be Riemannian manifolds and let

$$\pi:M\to B$$

be a smooth and surjective map

• This is a Riemannian submersion when $d\pi$ is surjective and

$$d\pi: (\ker d\pi)^{\perp} o \mathscr{X}(B) \otimes_{C^{\infty}(B)} C^{\infty}(M)$$

is an isometric isomorphism

This gives rise to vertical and horizontal vector fields

$$\mathscr{X}(M) \cong \mathscr{X}_V(M) \oplus \mathscr{X}_H(M).$$

On $\mathscr{X}(M)$ one can introduce the direct sum connection

$$\nabla^{\oplus} = \nabla^{M}_{V} \oplus \pi^* \nabla^{B},$$

and associate to it the second fundamental form and the curvature Ω of the fibration which yield a tensor $\omega \in \Omega^1(M) \otimes_{C^{\infty}(M)} \Omega^2(M)$

Proposition (Bismut, 1986)

The Levi-Civita connection ∇^M is related to the direct sum connection ∇^{\oplus} by the following formula

$$\langle \nabla_X^M Y, Z \rangle_M = \langle \nabla_X^{\oplus} Y, Z \rangle_M + \omega(X)(Y, Z).$$

Spin geometry and Clifford modules

Suppose that M and B are (even-dim) spin^c manifolds, so we have

$$Cl(M) \cong End_{C^{\infty}(M)}(\mathscr{E}_{M}), \qquad Cl(B) \cong End_{C^{\infty}(B)}(\mathscr{E}_{B})$$

and hermitian Clifford connections $\nabla^{\mathscr{E}_{M}}$ and $\nabla^{\mathscr{E}_{B}}$.

• We define the horizontal spinor module:

$$\mathscr{E}_H := \mathscr{E}_B \otimes_{C^{\infty}(B)} C^{\infty}(M); \qquad \nabla^{\mathscr{E}_H} := \pi^* \nabla^{\mathscr{E}_B}$$
 and $\mathsf{Cl}_H(M) \cong \mathsf{End}_{C^{\infty}(M)}(\mathscr{E}_H)$

We define the vertical spinor module:

$$\mathscr{E}_V := \mathscr{E}_H^* \otimes_{\mathsf{Cl}_H(M)} \mathscr{E}_M; \qquad
abla_X^{\mathscr{E}_V} =
abla_X^{\mathscr{E}_H^*} \otimes 1 + 1 \otimes
abla_X^{\mathscr{E}_M} + rac{1}{4} c(\omega(X))$$
 and $\mathsf{Cl}_V(M) \cong \mathsf{End}_{C^\infty(M)}(\mathscr{E}_V)$

Finally,

$$\mathscr{E}_H \otimes_{C^{\infty}(M)} \mathscr{E}_V \cong \mathscr{E}_M,$$

compatibly with $Cl_H(M) \widehat{\otimes}_{C^{\infty}(M)} Cl_V(M) \cong Cl(M)$.

The vertical operator

• We thus have a C^* -correspondence X from $C_0(M)$ to $C_0(B)$ by completing $\mathscr{E}_V^c := \mathscr{E}_V \otimes_{C^{\infty}(M)} C_c^{\infty}(M)$ with respect to

$$\langle \phi_1, \phi_2 \rangle_{X}(b) := \int_{F_b} \langle \phi_1, \phi_2 \rangle_{\mathscr{E}_V^c}$$

The following local expression defines a symmetric operator

$$(D_V)_0(\xi) = i \sum_{i=1}^{\dim(F)} c(e_i) \nabla_{e_i}^{\mathcal{E}_V}(\xi); \qquad (\xi \in \mathcal{E}_V^c)$$

Proposition

If $\pi: M \to B$ is a proper Riemannian submersion of spin^c manifolds, then (X, D_V) is an even unbounded Kasparov module from $C_0(M)$ to $C_0(B)$.

The horizontal operator and the connection

• The Dirac operator D_B : Dom $(D_B) \to L^2(\mathscr{E}_B^c)$ is locally

$$D_B = i \sum_{\alpha=1}^{\dim(B)} c(f_\alpha) \nabla_{f_\alpha}^{\mathscr{E}_B^c} : \mathscr{E}_B^c \to L^2(\mathscr{E}_B^c)$$

and $(C_c^{\infty}(B), L^2(\mathcal{E}_B^2), D_B)$ is a half-closed chain [Hil, 2010]

• The following defines a hermitian connection on X

$$\nabla_Z^X(\xi) = \nabla_{Z_H}^{\mathscr{E}_V}(\xi) + \frac{1}{2}k(Z_H) \cdot \xi$$

with $k \in \Omega^1(M)$ the mean curvature.

Lemma

The following local expression defines an odd symmetric unbounded operator in $X \widehat{\otimes}_{C_0(B)} L^2(\mathscr{E}_B^c)$:

$$(1 \otimes_{\nabla} D_B)(\xi \otimes r) := \xi \otimes D_B r + i \sum_{\alpha} \nabla^X_{f_{\alpha}}(\xi) \otimes c(f_{\alpha}) r.$$

The tensor sum

The tensor sum we are after is given by

$$D_V \widehat{\otimes} \gamma_B + 1 \widehat{\otimes}_\nabla D_B : \mathsf{Dom}(D_V \times_\nabla D_B)_0 \to X \widehat{\otimes}_{C_0(B)} L^2(\mathscr{E}_B^c)$$

- The closure of this symmetric operator is denoted $D_V \times_{\nabla} D_B$.
- Since $P \subseteq M$ has codimension > 1, it follows that the closures of the following two operators coincide:
 - 1 the first-order differential operator $D_{\overline{M}}$ on $\mathscr{E}_{\overline{M}} := \Gamma^{\infty}(\overline{M}, \mathscr{E}_{\overline{M}})$
 - 2 the restriction of $D_{\overline{M}}$ to $\mathscr{E}_{M}^{c} := \Gamma_{c}^{\infty}(M, \mathscr{E}_{\overline{M}})$ where $M = \overline{M} \setminus \overline{P}$

Theorem (KvS, 2017)

Under the unitary isomorphism $V: X \widehat{\otimes}_{C_0(B)} L^2(\mathscr{E}_B^c) \to L^2(\mathscr{E}_{\overline{M}})$ we have the identity of selfadjoint operators

$$V(D_V \times_{\nabla} D_B)V^* = D_{\overline{M}} + \widetilde{c}(\Omega).$$

Moreover, the tensor sum is an unbounded representative of

$$KK(C_0(M), C_0(B)) \otimes_{C_0(B)} KK(C_0(B), \mathbb{C}) \to KK(C_0(M), \mathbb{C})$$

Application: toric noncommutative manifolds

- Consider the class of examples coming from actions of a torus G on a spin^c manifold \overline{M} with principal stratum M and with B = M/G a spin^c manifold.
- Prototype: the four-sphere S^4 with a \mathbb{T}^2 -action (the starting point for the Connes-Landi four-sphere)

The four-sphere

• Toroidal coordinates: $0 < \theta_1, \theta_2 < 2\pi, 0 < \varphi < \pi/2$ $-\pi/2 \le \psi \le \pi/2$, and write

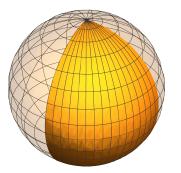
$$z_1 = e^{i\theta_1} \cos \varphi \cos \psi;$$

 $z_2 = e^{i\theta_2} \sin \varphi \cos \psi;$
 $x = \sin \psi$

• Action of \mathbb{T}^2 is by translating the θ_1, θ_2 -coordinates

• The orbit space $\mathbb{S}^4/\mathbb{T}^2\cong Q^2$ is a closed quadrant in the two-sphere, parametrized by

$$0 \le \varphi \le \pi/2, \qquad -\pi/2 \le \psi \le \pi/2$$



- Principal stratum \mathbb{S}_0^4 is a trivial \mathbb{T}^2 -principal fiber bundle over the interior Q_0^2 of Q^2
- Moreover, $\pi: \mathbb{S}_0^4 \to Q_0^2$ is a Riemannian submersion for the metric on Q_0^2 induced by the round metric on \mathbb{S}^2

Spin geometry of \mathbb{S}^4

• The Dirac operator for the round metric is a \mathbb{T}^2 -invariant selfadjoint operator

$$D_{\mathbb{S}^4}: \mathsf{Dom}(D_{\mathbb{S}^4}) o L^2(\mathbb{S}^4, \mathscr{E}_{\mathbb{S}^4})$$

with local expression:

$$\begin{split} D_{\mathbb{S}^4} &= i \frac{1}{\cos \varphi \cos \psi} \gamma^1 \frac{\partial}{\partial \theta_1} + i \frac{1}{\sin \varphi \cos \psi} \gamma^2 \frac{\partial}{\partial \theta_2} \\ &+ i \frac{1}{\cos \psi} \gamma^3 \left(\frac{\partial}{\partial \varphi} + \frac{1}{2} \cot \varphi - \frac{1}{2} \tan \varphi \right) + i \gamma^4 \left(\frac{\partial}{\partial \psi} - \frac{3}{2} \tan \psi \right) \end{split}$$

 Since the 'subprincipal' stratum is of codimension two, it follows that $C_c^{\infty}(\mathbb{S}_0^4) \otimes \mathbb{C}^4 \subset L^2(\mathbb{S}^4, \mathcal{E}_{\mathbb{S}^4})$ is a core for $D_{\mathbb{S}^4}$

Vertical operator

• We consider a C^* -correspondence X from $C_0(\mathbb{S}_0^4)$ to $C_0(\mathbb{Q}_0^2)$, defined as the Hilbert C^* -completion of $C_c^{\infty}(\mathbb{S}_0^4) \otimes \mathbb{C}^2$ wrt

$$\langle \mathbf{s}, \mathbf{t} \rangle_{X} = \int_{\mathbb{T}^{2}} \overline{\mathbf{s}(\theta_{1}, \theta_{2}, \varphi, \psi)} \cdot \mathbf{t}(\theta_{1}, \theta_{2}, \varphi, \psi) d\theta_{1} d\theta_{2} \cdot \sin \varphi \cos \varphi \cos^{2} \psi,$$

There is a unitary isomorphism (induced by pointw. multipl.)

$$u: X \otimes_{C_0(Q_0^2)} L^2(Q_0^2) \otimes \mathbb{C}^2 \to L^2(\mathbb{S}_0^4) \otimes \mathbb{C}^4$$

• We define a symmetric operator $(D_V)_0: C_c^{\infty}(\mathbb{S}_0^4) \otimes \mathbb{C}^2 \to X$:

$$(D_V)_0 = i \frac{1}{\cos \varphi \cos \psi} \sigma^1 \frac{\partial}{\partial \theta_1} + i \frac{1}{\sin \varphi \cos \psi} \sigma^2 \frac{\partial}{\partial \theta_2}.$$

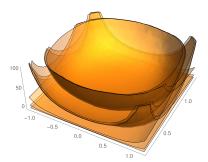
We denote its closure by D_V : Dom $(D_V) \to X$.

• One can find families of eigenvectors $\Psi_{n_1n_2}^{\pm}$ in $C^{\infty}(\mathbb{T}^2)\otimes\mathbb{C}^2$.

$$D_V\left(f\Psi_{n1,n_2}^{\pm}\right)=\pm\lambda_{n_1n_2}\cdot f\Psi_{n_1n_2}^{\pm}.$$

for any $f \in C_c^{\infty}(Q_0^2)$ where

$$\lambda_{n_1 n_2} = \sqrt{\frac{n_1^2}{\cos^2 \varphi \cos^2 \psi} + \frac{n_2^2}{\sin^2 \varphi \cos^2 \psi}}$$

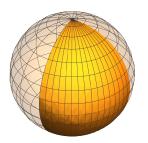


• Regularity, selfadjointness and compactness of the resolvent: D_V defines unbdd Kasparov module from $C_0(\mathbb{S}_0^4)$ to $C_0(\mathbb{Q}_0^2)$

Horizontal operator

We define an operator $(D_{Q_0^2})_0: C_c^{\infty}(Q_0^2) \otimes \mathbb{C}^2 \to L^2(Q_0^2) \otimes \mathbb{C}^2$ as the restriction of the Dirac operator on \mathbb{S}^2 to the open quadrant Q_0^2 :

$$(D_{Q_0^2})_0 = i \frac{1}{\cos \psi} \sigma^1 \frac{\partial}{\partial \varphi} + i \sigma^2 \left(\frac{\partial}{\partial \psi} - \frac{1}{2} \tan \psi \right).$$



This is a symmetric operator. Its closure $D_{Q_0^2}$ defines a half-closed chain [Hil, 2010] from $C_0(Q_0^2)$ to $\mathbb C$

Tensor sum of D_V and $D_{Q_0^2}$

We lift the operators D_V and $D_{Q_0^2}$ to $X \otimes_{C_0(Q_0^2)} L^2(Q_0^2) \otimes \mathbb{C}^2$.

• For D_V we take $D_V \widehat{\otimes} \gamma$ and have on $C_c^{\infty}(\mathbb{S}_0^4) \otimes \mathbb{C}^4$ that

$$u(D_V\widehat{\otimes}\gamma)u^* = i\frac{1}{\cos\varphi\cos\psi}\gamma^1\frac{\partial}{\partial\theta_1} + i\frac{1}{\sin\varphi\cos\psi}\gamma^2\frac{\partial}{\partial\theta_2},$$

• For the operator $D_{Q_0^2}$, we first need a hermitian connection on X

$$\nabla_{\partial/\partial\varphi} = \frac{\partial}{\partial\varphi} + \frac{1}{2}\cot\varphi - \frac{1}{2}\tan\varphi \qquad \nabla_{\partial/\partial\psi} = \frac{\partial}{\partial\psi} - \tan\psi$$

• The operator $1 \widehat{\otimes}_{\nabla} D_{Q_0^2}$ then becomes

$$\begin{split} u(1\widehat{\otimes}_{\nabla}D_{Q_0^2})u^* &= i\frac{1}{\cos\psi}\gamma^3\left(\frac{\partial}{\partial\varphi} + \frac{1}{2}\cot\varphi - \frac{1}{2}\tan\varphi\right) \\ &+ i\gamma^4\left(\frac{\partial}{\partial\psi} - \frac{3}{2}\tan\psi\right) \end{split}$$

• The tensor sum we are after is given by

$$D_{V}\widehat{\otimes}\gamma + 1\widehat{\otimes}_{\nabla}D_{Q_{0}^{2}}$$

• This is a symmetric operator and we denote its closure by

$$D_V imes_{
abla} D_{Q_0^2} : \mathsf{Dom}(D_V imes_{
abla} D_{Q_0^2}) o X \otimes_{C_0(Q_0^2)} L^2(Q_0^2) \otimes \mathbb{C}^2.$$

• Since $C_c^{\infty}(\mathbb{S}_0^4) \otimes \mathbb{C}^2 \otimes_{C_c^{\infty}(Q_0^2)} C_c^{\infty}(Q_0^2) \otimes \mathbb{C}^2$ is a core for $D_V \times_{\nabla} D_{Q_0^2}$ which is mapped by u to the core $C_c^{\infty}(\mathbb{S}_0^4) \otimes \mathbb{C}^4$ of $D_{\mathbb{S}^4}$, it follows that

$$u(D_V \times_{\nabla} D_{Q_0^2})u^* = D_{\mathbb{S}^4}$$

as an equality on $\mathsf{Dom}(D_{\mathbb{S}^4})$

 A connection and positivity property imply that this is an unbounded representative of the internal product

$$KK(C_0(\mathbb{S}^4_0), C_0(Q_0^2)) \otimes_{C_0(Q_0^2)} KK(C_0(Q_0^2), \mathbb{C}) \to KK(C_0(\mathbb{S}^4_0), \mathbb{C})$$

Summary

For almost-regular fibrations (M, P, B, π) of spin^c manifolds we have obtained with $M = \overline{M} \setminus \overline{P}$:

- an unbounded Kasparov module (X, D_V) from $C_0(M)$ to $C_0(B)$ describing the vertical data
- a connection ∇ on X that allows to lift the Dirac operator D_B to a symmetric operator $1 \widehat{\otimes}_{\nabla} D_B$ on $X \widehat{\otimes}_{C_0(B)} L^2(\mathscr{E}_B^c)$
- an identity of selfadjoint operators on $Dom(D_{\overline{M}})$

$$V\left(D_V\widehat{\otimes}\gamma+1\widehat{\otimes}_{\nabla}D_B\right)V^*=D_{\overline{M}}+\widetilde{c}(\Omega).$$

 A local connection and positivity property imply that this is an unbounded representative of the internal product

$$\mathsf{KK}(\mathsf{C}_0(M),\mathsf{C}_0(B))\otimes_{\mathsf{C}_0(B)}\mathsf{KK}(\mathsf{C}_0(M),\mathbb{C})\to\mathsf{KK}(\mathsf{C}_0(M),\mathbb{C})$$

 Applications: tori acting on compact Riemannian spin^c manifolds

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